

The gamma and the beta function

As mentioned in the book [1], see page 6, the integral representation (1.1.18) is often taken as a definition for the gamma function $\Gamma(z)$. The advantage of this alternative definition is that we might avoid the use of infinite products (see appendix A).

Definition 1.

$$\Gamma(z) = \int_0^{\infty} e^{-t} t^{z-1} dt, \quad \operatorname{Re} z > 0. \quad (1)$$

From this definition it is clear that $\Gamma(z)$ is analytic for $\operatorname{Re} z > 0$. By using integration by parts we find that

$$\begin{aligned} \Gamma(z+1) &= \int_0^{\infty} e^{-t} t^z dt = - \int_0^{\infty} t^z de^{-t} = -e^{-t} t^z \Big|_0^{\infty} + \int_0^{\infty} e^{-t} dt^z \\ &= z \int_0^{\infty} e^{-t} t^{z-1} dt = z\Gamma(z), \quad \operatorname{Re} z > 0. \end{aligned}$$

Hence we have

Theorem 1.

$$\Gamma(z+1) = z\Gamma(z), \quad \operatorname{Re} z > 0. \quad (2)$$

Further we have

$$\Gamma(1) = \int_0^{\infty} e^{-t} dt = -e^{-t} \Big|_0^{\infty} = 1. \quad (3)$$

Combining (2) and (3), this leads to

$$\Gamma(n+1) = n!, \quad n = 0, 1, 2, \dots \quad (4)$$

The functional relation (2) can be used to find an analytic continuation of the gamma function for $\operatorname{Re} z \leq 0$. For $\operatorname{Re} z > 0$ the gamma function $\Gamma(z)$ is defined by (1). The functional relation (2) also holds for $\operatorname{Re} z > 0$.

Let $-1 < \operatorname{Re} z \leq 0$, then we have $\operatorname{Re}(z+1) > 0$. Hence, $\Gamma(z+1)$ is defined by the integral representation (1). Now we define

$$\Gamma(z) = \frac{\Gamma(z+1)}{z}, \quad -1 < \operatorname{Re} z \leq 0, \quad z \neq 0.$$

Then the gamma function $\Gamma(z)$ is analytic for $\operatorname{Re} z > -1$ except $z = 0$. For $z = 0$ we have

$$\lim_{z \rightarrow 0} z\Gamma(z) = \lim_{z \rightarrow 0} \Gamma(z+1) = \Gamma(1) = 1.$$

This implies that $\Gamma(z)$ has a single pole at $z = 0$ with residue 1.

This process can be repeated for $-2 < \operatorname{Re} z \leq -1$, $-3 < \operatorname{Re} z \leq -2$, etcetera. Then the gamma function turns out to be an analytic function on \mathbb{C} except for single poles at $z = 0, -1, -2, \dots$. The residue at $z = -n$ equals

$$\begin{aligned} \lim_{z \rightarrow -n} (z+n)\Gamma(z) &= \lim_{z \rightarrow -n} (z+n) \frac{\Gamma(z+1)}{z} = \lim_{z \rightarrow -n} (z+n) \frac{1}{z} \frac{1}{z+1} \cdots \frac{1}{z+n-1} \frac{\Gamma(z+n+1)}{z+n} \\ &= \frac{\Gamma(1)}{(-n)(-n+1)\cdots(-1)} = \frac{(-1)^n}{n!}, \quad n = 0, 1, 2, \dots \end{aligned}$$

As indicated in the book [1], see page 8, the limit formula (1.1.5) can be obtained from the integral representation (1) by using induction as follows. We first prove that

$$\int_0^1 (1-t)^n t^{z-1} dt = \frac{n!}{(z)_{n+1}} \quad (5)$$

for $\operatorname{Re} z > 0$ and $n = 0, 1, 2, \dots$. Here the shifted factorial $(a)_k$ is defined by

Definition 2.

$$(a)_k = a(a+1)\cdots(a+k-1), \quad k = 1, 2, 3, \dots \quad \text{and} \quad (a)_0 = 1. \quad (6)$$

In order to prove (5) by induction we first take $n = 0$ to obtain for $\operatorname{Re} z > 0$

$$\int_0^1 t^{z-1} dt = \left. \frac{t^z}{z} \right|_0^1 = \frac{1}{z} = \frac{0!}{(z)_1}.$$

Now we assume that (5) holds for $n = k$. Then we have

$$\begin{aligned} \int_0^1 (1-t)^{k+1} t^{z-1} dt &= \int_0^1 (1-t)(1-t)^k t^{z-1} dt = \int_0^1 (1-t)^k t^{z-1} dt - \int_0^1 (1-t)^k t^z dt \\ &= \frac{k!}{(z)_{k+1}} - \frac{k!}{(z+1)_{k+1}} = \frac{k!}{(z)_{k+2}}(z+k+1-z) = \frac{(k+1)!}{(z)_{k+2}}, \end{aligned}$$

which is (5) for $n = k + 1$. This proves that (5) holds for all $n = 0, 1, 2, \dots$

Now we set $t = u/n$ into (5) to find that

$$\frac{1}{n^z} \int_0^n \left(1 - \frac{u}{n}\right)^n u^{z-1} du = \frac{n!}{(z)_{n+1}} \implies \int_0^n \left(1 - \frac{u}{n}\right)^n u^{z-1} du = \frac{n! n^z}{(z)_{n+1}}.$$

Since we have

$$\lim_{n \rightarrow \infty} \left(1 - \frac{u}{n}\right)^n = e^{-u},$$

we conclude that

$$\Gamma(z) = \int_0^\infty e^{-u} u^{z-1} du = \lim_{n \rightarrow \infty} \frac{n! n^z}{(z)_{n+1}}.$$

The beta function $B(u, v)$ is also defined by means of an integral:

Definition 3.

$$B(u, v) = \int_0^1 t^{u-1} (1-t)^{v-1} dt, \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0. \quad (7)$$

This integral is often called the beta integral. From the definition we easily obtain the symmetry

$$B(u, v) = B(v, u), \quad (8)$$

since we have by using the substitution $t = 1 - s$

$$B(u, v) = \int_0^1 t^{u-1} (1-t)^{v-1} dt = - \int_1^0 (1-s)^{u-1} s^{v-1} ds = \int_0^1 s^{v-1} (1-s)^{u-1} ds = B(v, u).$$

The connection between the beta function and the gamma function is given by the following theorem:

Theorem 2.

$$B(u, v) = \frac{\Gamma(u)\Gamma(v)}{\Gamma(u+v)}, \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0. \quad (9)$$

In order to prove this theorem we use the definition (1) to obtain

$$\Gamma(u)\Gamma(v) = \int_0^\infty e^{-t} t^{u-1} dt \int_0^\infty e^{-s} s^{v-1} ds = \int_0^\infty \int_0^\infty e^{-(t+s)} t^{u-1} s^{v-1} dt ds.$$

Now we apply the change of variables $t = xy$ and $s = x(1 - y)$ to this double integral. Note that $t + s = x$ and that $0 < t < \infty$ and $0 < s < \infty$ imply that $0 < x < \infty$ and $0 < y < 1$. The Jacobian of this transformation is

$$\frac{\partial(t, s)}{\partial(x, y)} = \begin{vmatrix} y & x \\ 1 - y & -x \end{vmatrix} = -xy - x + xy = -x.$$

Since $x > 0$ we conclude that $dt ds = \left| \frac{\partial(t, s)}{\partial(x, y)} \right| dx dy = x dx dy$. Hence we have

$$\begin{aligned} \Gamma(u)\Gamma(v) &= \int_0^1 \int_0^\infty e^{-x} x^{u-1} y^{u-1} x^{v-1} (1-y)^{v-1} x dx dy \\ &= \int_0^\infty e^{-x} x^{u+v-1} dx \int_0^1 y^{u-1} (1-y)^{v-1} dy = \Gamma(u+v)B(u, v). \end{aligned}$$

This proves (9).

There exist many useful forms of the beta integral which can be obtained by an appropriate change of variables. For instance, if we set $t = s/(s+1)$ into (7) we obtain

$$\begin{aligned} B(u, v) &= \int_0^1 t^{u-1} (1-t)^{v-1} dt = \int_0^\infty s^{u-1} (s+1)^{-u+1} (s+1)^{-v+1} (s+1)^{-2} ds \\ &= \int_0^\infty \frac{s^{u-1}}{(s+1)^{u+v}} ds, \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0. \end{aligned}$$

This proves

Theorem 3.

$$B(u, v) = \int_0^\infty \frac{s^{u-1}}{(s+1)^{u+v}} ds, \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0. \quad (10)$$

If we set $t = \cos^2 \varphi$ into (7) we find that

$$\begin{aligned} B(u, v) &= \int_0^1 t^{u-1} (1-t)^{v-1} dt = -2 \int_{\pi/2}^0 (\cos \varphi)^{2u-2} (\sin \varphi)^{2v-2} \cos \varphi \sin \varphi d\varphi \\ &= 2 \int_0^{\pi/2} (\cos \varphi)^{2u-1} (\sin \varphi)^{2v-1} d\varphi, \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0. \end{aligned}$$

Hence we have

Theorem 4.

$$B(u, v) = 2 \int_0^{\pi/2} (\cos \varphi)^{2u-1} (\sin \varphi)^{2v-1} d\varphi, \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0. \quad (11)$$

Finally, the substitution $t = (s - a)/(b - a)$ in (7) leads to

$$\begin{aligned} B(u, v) &= \int_0^1 t^{u-1}(1-t)^{v-1} dt \\ &= \int_a^b (s-a)^{u-1}(b-a)^{-u+1}(b-s)^{v-1}(b-a)^{-v+1}(b-a)^{-1} ds \\ &= (b-a)^{-u-v+1} \int_a^b (s-a)^{u-1}(b-s)^{v-1} ds, \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0. \end{aligned}$$

Hence we have

Theorem 5.

$$\int_a^b (s-a)^{u-1}(b-s)^{v-1} ds = (b-a)^{u+v-1} B(u, v), \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0. \quad (12)$$

The special case $a = -1$ and $b = 1$ is of special interest as we will see later:

$$\int_{-1}^1 (1+s)^{u-1}(1-s)^{v-1} ds = 2^{u+v-1} B(u, v), \quad \operatorname{Re} u > 0, \quad \operatorname{Re} v > 0.$$

The different forms for the beta function have a lot of consequences. For instance, if we set $u = v = 1/2$ in (9) we find that

$$B(1/2, 1/2) = \frac{\Gamma(1/2)\Gamma(1/2)}{\Gamma(1)} = \{\Gamma(1/2)\}^2.$$

On the other hand, we have by using (11)

$$B(1/2, 1/2) = 2 \int_0^{\pi/2} d\varphi = 2 \cdot \frac{\pi}{2} = \pi.$$

This implies that

$$\Gamma(1/2) = \sqrt{\pi}. \quad (13)$$

By using the transformation $x^2 = t$ we now easily obtain the value of the normal integral

$$\int_{-\infty}^{\infty} e^{-x^2} dx = 2 \int_0^{\infty} e^{-x^2} dx = \int_0^{\infty} e^{-t} t^{-1/2} dt = \Gamma(1/2) = \sqrt{\pi}. \quad (14)$$

The combination of (9) and (11) can be used to compute integrals such as

$$\int_0^{\pi/2} (\cos \varphi)^5 (\sin \varphi)^7 d\varphi = \frac{1}{2} \cdot B(3, 4) = \frac{1}{2} \cdot \frac{\Gamma(3)\Gamma(4)}{\Gamma(7)} = \frac{1}{2} \cdot \frac{2!3!}{6!} = \frac{1}{2} \cdot \frac{2}{4 \cdot 5 \cdot 6} = \frac{1}{120},$$

$$\begin{aligned} \int_0^{\pi/2} (\cos \varphi)^7 (\sin \varphi)^4 d\varphi &= \frac{1}{2} \cdot B(4, 5/2) = \frac{1}{2} \cdot \frac{\Gamma(4)\Gamma(5/2)}{\Gamma(13/2)} = \frac{1}{2} \cdot \frac{3!}{\frac{5}{2} \cdot \frac{7}{2} \cdot \frac{9}{2} \cdot \frac{11}{2}} \\ &= \frac{1}{2} \cdot \frac{6 \cdot 2^4}{5 \cdot 7 \cdot 9 \cdot 11} = \frac{16}{1155} \end{aligned}$$

and

$$\begin{aligned} \int_0^{\pi/2} (\cos \varphi)^4 (\sin \varphi)^6 d\varphi &= \frac{1}{2} \cdot B(5/2, 7/2) = \frac{1}{2} \cdot \frac{\Gamma(5/2)\Gamma(7/2)}{\Gamma(6)} \\ &= \frac{1}{2} \cdot \frac{\frac{3}{2} \cdot \frac{1}{2} \cdot \Gamma(1/2) \cdot \frac{5}{2} \cdot \frac{3}{2} \cdot \frac{1}{2} \cdot \Gamma(1/2)}{5!} = \frac{5 \cdot 3^2 \cdot \pi}{2^6 \cdot 2 \cdot 3 \cdot 4 \cdot 5} = \frac{3\pi}{2^9} = \frac{3\pi}{512}. \end{aligned}$$

Another important consequence of (9) and (11) is Legendre's duplication formula for the gamma function:

Theorem 6.

$$\Gamma(z)\Gamma(z + 1/2) = 2^{1-2z} \sqrt{\pi} \Gamma(2z), \quad \operatorname{Re} z > 0. \quad (15)$$

In order to prove this we use (11) and the transformation $2\varphi = \tau$ to find that

$$\begin{aligned} B(z, z) &= 2 \int_0^{\pi/2} (\cos \varphi)^{2z-1} (\sin \varphi)^{2z-1} d\varphi = 2 \cdot 2^{1-2z} \int_0^{\pi/2} (\sin 2\varphi)^{2z-1} d\varphi \\ &= 2^{1-2z} \int_0^{\pi} (\sin \tau)^{2z-1} d\tau = 2^{1-2z} \cdot 2 \int_0^{\pi/2} (\sin \tau)^{2z-1} d\tau = 2^{1-2z} \cdot B(z, 1/2). \end{aligned}$$

Now we apply (9) to obtain

$$\frac{\Gamma(z)\Gamma(z)}{\Gamma(2z)} = B(z, z) = 2^{1-2z} \cdot B(z, 1/2) = 2^{1-2z} \cdot \frac{\Gamma(z)\Gamma(1/2)}{\Gamma(z + 1/2)}, \quad \operatorname{Re} z > 0.$$

Finally, by using (13), this implies that

$$\Gamma(z)\Gamma(z + 1/2) = 2^{1-2z} \sqrt{\pi} \Gamma(2z), \quad \operatorname{Re} z > 0.$$

This proves the theorem.

Legendre's duplication formula can be generalized to Gauss's multiplication formula:

Theorem 7.

$$\Gamma(z) \prod_{k=1}^{n-1} \Gamma(z + k/n) = n^{1/2-nz} (2\pi)^{(n-1)/2} \Gamma(nz), \quad n \in \{1, 2, 3, \dots\}. \quad (16)$$

The case $n = 1$ is trivial and the case $n = 2$ is Legendre's duplication formula.

Another property of the gamma function is given by Euler's reflection formula:

Theorem 8.

$$\Gamma(z)\Gamma(1-z) = \frac{\pi}{\sin \pi z}, \quad z \neq 0, \pm 1, \pm 2, \dots \quad (17)$$

This can be shown by using contour integration in the complex plane as follows. First we restrict to real values of z , say $z = x$ with $0 < x < 1$. By using (9) and (10) we have

$$\Gamma(x)\Gamma(1-x) = B(x, 1-x) = \int_0^{\infty} \frac{t^{x-1}}{t+1} dt.$$

In order to compute this integral we consider the contour integral

$$\int_{\mathcal{C}} \frac{z^{x-1}}{1-z} dz,$$

where the contour \mathcal{C} consists of two circles about the origin of radii R and ϵ respectively, which are joined along the negative real axis from $-R$ to $-\epsilon$. Move along the outer circle with radius R in the positive (counterclockwise) direction and along the inner circle with radius ϵ in the negative (clockwise) direction. By the residue theorem we have

$$\int_{\mathcal{C}} \frac{z^{x-1}}{1-z} dz = -2\pi i,$$

when z^{x-1} has its principal value. This implies that

$$-2\pi i = \int_{\mathcal{C}_1} \frac{z^{x-1}}{1-z} dz + \int_{\mathcal{C}_2} \frac{z^{x-1}}{1-z} dz + \int_{\mathcal{C}_3} \frac{z^{x-1}}{1-z} dz + \int_{\mathcal{C}_4} \frac{z^{x-1}}{1-z} dz,$$

where \mathcal{C}_1 denotes the outer circle with radius R , \mathcal{C}_2 denotes the line segment from $-R$ to $-\epsilon$, \mathcal{C}_3 denotes the inner circle with radius ϵ and \mathcal{C}_4 denotes the line segment from $-\epsilon$ to $-R$. Then we have by writing $z = Re^{i\theta}$ for the outer circle

$$\int_{\mathcal{C}_1} \frac{z^{x-1}}{1-z} dz = \int_{-\pi}^{\pi} \frac{R^{x-1} e^{i(x-1)\theta}}{1 - Re^{i\theta}} d(Re^{i\theta}) = \int_{-\pi}^{\pi} \frac{iR^x e^{ix\theta}}{1 - Re^{i\theta}} d\theta.$$

In the same way we have by writing $z = \epsilon e^{i\theta}$ for the inner circle

$$\int_{\mathcal{C}_3} \frac{z^{x-1}}{1-z} dz = \int_{\pi}^{-\pi} \frac{i\epsilon^x e^{ix\theta}}{1 - \epsilon e^{i\theta}} d\theta.$$

For the line segment from $-R$ to $-\epsilon$ we have by writing $z = -t = te^{\pi i}$

$$\int_{\mathcal{C}_2} \frac{z^{x-1}}{1-z} dz = \int_R^{\epsilon} \frac{t^{x-1} e^{i(x-1)\pi}}{1+t} d(te^{\pi i}) = \int_R^{\epsilon} \frac{t^{x-1} e^{ix\pi}}{1+t} dt.$$

In the same way we have by writing $z = -t = te^{-\pi i}$

$$\int_{\mathcal{C}_4} \frac{z^{x-1}}{1-z} dz = \int_{\epsilon}^R \frac{t^{x-1} e^{-ix\pi}}{1+t} dt.$$

Since $0 < x < 1$ we have

$$\lim_{R \rightarrow \infty} \int_{-\pi}^{\pi} \frac{iR^x e^{ix\theta}}{1 - Re^{i\theta}} d\theta = 0 \quad \text{and} \quad \lim_{\epsilon \downarrow 0} \int_{\pi}^{-\pi} \frac{i\epsilon^x e^{ix\theta}}{1 - \epsilon e^{i\theta}} d\theta = 0.$$

Hence we have

$$-2\pi i = \int_{\infty}^0 \frac{t^{x-1} e^{ix\pi}}{1+t} dt + \int_0^{\infty} \frac{t^{x-1} e^{-ix\pi}}{1+t} dt,$$

or

$$-2\pi i = (e^{-ix\pi} - e^{ix\pi}) \int_0^{\infty} \frac{t^{x-1}}{1+t} dt \implies \int_0^{\infty} \frac{t^{x-1}}{1+t} dt = \frac{2\pi i}{e^{ix\pi} - e^{-ix\pi}} = \frac{\pi}{\sin \pi x}.$$

This proves the theorem for real values of z , say $z = x$ with $0 < x < 1$. The full result follows by analytic continuation. Alternatively, the result can be obtained as follows. If (17) holds for real values of z with $0 < z < 1$, then it holds for all complex z with $0 < \operatorname{Re} z < 1$ by

analyticity. Then it also holds for $\operatorname{Re} z = 0$ with $z \neq 0$ by continuity. Finally, the full result follows for z shifted by integers using (2) and $\sin(z + \pi) = -\sin z$. Note that (17) holds for all complex values of z with $z \neq 0, -1, -2, \dots$. Instead of (17) we may write

$$\frac{1}{\Gamma(z)\Gamma(1-z)} = \frac{\sin \pi z}{\pi}, \quad (18)$$

which holds for all $z \in \mathbb{C}$.

Now we will prove an asymptotic formula which is due to Stirling. First we define

Definition 4. *Two functions f and g of a real variable x are called asymptotically equal, notation*

$$f \sim g \quad \text{for } x \rightarrow \infty, \quad \text{if } \lim_{x \rightarrow \infty} \frac{f(x)}{g(x)} = 1.$$

Now we have Stirling's formula:

Theorem 9.

$$\Gamma(x+1) \sim x^{x+1/2} e^{-x} \sqrt{2\pi}, \quad x \rightarrow \infty. \quad (19)$$

Here x denotes a real variable. This can be proved as follows. Consider

$$\Gamma(x+1) = \int_0^\infty e^{-t} t^x dt,$$

where $x \in \mathbb{R}$. Then we obtain by using the transformation $t = x(1+u)$

$$\begin{aligned} \Gamma(x+1) &= \int_{-1}^\infty e^{-x(1+u)} x^x (1+u)^x x du = x^{x+1} e^{-x} \int_{-1}^\infty e^{-xu} (1+u)^x du \\ &= x^{x+1} e^{-x} \int_{-1}^\infty e^{x(-u+\ln(1+u))} du. \end{aligned}$$

The function $f(u) = -u + \ln(1+u)$ equals zero for $u = 0$. For other values of u we have $f(u) < 0$. This implies that the integrand of the last integral equals 1 at $u = 0$ and that this integrand becomes very small for large values of x at other values of u . So for large values of x we only have to deal with the integrand near $u = 0$. Note that we have

$$f(u) = -u + \ln(1+u) = -\frac{1}{2}u^2 + \mathcal{O}(u^3) \quad \text{for } u \rightarrow 0.$$

This implies that

$$\int_{-1}^\infty e^{x(-u+\ln(1+u))} du \sim \int_{-\infty}^\infty e^{-xu^2/2} du \quad \text{for } x \rightarrow \infty.$$

If we set $u = t\sqrt{2/x}$ we have by using the normal integral (14)

$$\int_{-\infty}^\infty e^{-xu^2/2} du = x^{-1/2} \sqrt{2} \int_{-\infty}^\infty e^{-t^2} dt = x^{-1/2} \sqrt{2\pi}.$$

Hence we have

$$\Gamma(x+1) \sim x^{x+1/2} e^{-x} \sqrt{2\pi}, \quad x \rightarrow \infty,$$

which proves the theorem.

Note that Stirling's formula implies that

$$n! \sim n^n e^{-n} \sqrt{2\pi n} \quad \text{for } n \rightarrow \infty$$

and that

$$\frac{\Gamma(n+a)}{\Gamma(n+b)} \sim n^{a-b} \quad \text{for } n \rightarrow \infty.$$

The theorem can be extended for z in the complex plane:

Theorem 10. *For $\delta > 0$ we have*

$$\Gamma(z+1) \sim z^{z+1/2} e^{-z} \sqrt{2\pi} \quad \text{for } |z| \rightarrow \infty \quad \text{with } |\arg z| \leq \pi - \delta. \quad (20)$$

Stirling's asymptotic formula can be used to give an alternative proof for Euler's reflection formula (17) for the gamma function. Consider the function

$$f(z) = \Gamma(z)\Gamma(1-z)\sin\pi z.$$

Then we have

$$f(z+1) = \Gamma(z+1)\Gamma(-z)\sin\pi(z+1) = z\Gamma(z) \cdot \frac{\Gamma(1-z)}{-z} \cdot -\sin\pi z = f(z).$$

Hence, f is periodic with period 1. Further we have

$$\lim_{z \rightarrow 0} f(z) = \lim_{z \rightarrow 0} \Gamma(z)\Gamma(1-z)\sin\pi z = \lim_{z \rightarrow 0} \Gamma(z+1)\Gamma(1-z) \frac{\sin\pi z}{z} = \pi, \quad (21)$$

which implies that f has no poles. Hence, f is analytic and periodic with period 1. Now we want to apply Liouville's theorem for entire functions, id est functions which are analytic on the whole complex plane:

Theorem 11. *Every bounded entire function is constant.*

Therefore, we want to show that f is bounded. Since f is periodic with period 1 we consider $0 \leq \operatorname{Re} z \leq 1$, say $z = x + iy$ with x and y real and $0 \leq x \leq 1$. Then we have

$$\sin\pi z = \frac{e^{i\pi z} - e^{-i\pi z}}{2i} \sim -\frac{1}{2i} e^{-i\pi z} = \frac{i}{2} e^{-i\pi z} \quad \text{for } y \rightarrow \infty.$$

Now we apply Stirling's formula to obtain

$$f(z) = \Gamma(z)\Gamma(1-z)\sin\pi z \sim z^{z-1/2} e^{-z} \sqrt{2\pi} (-z)^{-z+1/2} e^z \sqrt{2\pi} \frac{i}{2} e^{-i\pi z}.$$

For $y > 0$ we have $-z/z = e^{-\pi i}$. Hence, $f(z) \sim \pi$ for $y \rightarrow \infty$. This implies that f is bounded. So, Liouville's theorem implies that f is constant. By using (21) we conclude that $f(z) = \pi$, which proves Euler's reflection formula (17) or (18).

Stirling's formula can also be used to give an alternative proof for Legendre's duplication formula (15). Consider the function

$$g(z) = 2^{2z-1} \frac{\Gamma(z)\Gamma(z+1/2)}{\Gamma(1/2)\Gamma(2z)}.$$

Then we have by using (2)

$$g(z+1) = 2^{2z+1} \frac{\Gamma(z+1)\Gamma(z+3/2)}{\Gamma(1/2)\Gamma(2z+2)} = 2^{2z+1} \frac{z\Gamma(z)(z+1/2)\Gamma(z+1/2)}{\Gamma(1/2)(2z+1)2z\Gamma(2z)} = g(z).$$

Further we have by using (13) and Stirling's asymptotic formula (20)

$$g(z) \sim 2^{2z-1} \frac{z^{z-1/2} e^{-z} \sqrt{2\pi} z^z e^{-z} \sqrt{2\pi}}{\sqrt{\pi} 2^{2z-1/2} z^{2z-1/2} e^{-2z} \sqrt{2\pi}} = 1.$$

This implies that

$$\lim_{n \rightarrow \infty} g(z+n) = 1,$$

also for integer values for n . On the other hand we have $g(z+n) = g(z)$ for integer values of n . This implies that $g(z) = 1$ for all z . This proves (15).

Finally we have the digamma function $\psi(z)$ which is related to the gamma function. This function $\psi(z)$ is defined as follows.

Definition 5.

$$\psi(z) = \frac{\Gamma'(z)}{\Gamma(z)} = \frac{d}{dz} \ln \Gamma(z), \quad z \neq 0, -1, -2, \dots \quad (22)$$

A property of this digamma function that is easily proved by using (2) is given by the following theorem:

Theorem 12.

$$\psi(z+1) = \psi(z) + \frac{1}{z}. \quad (23)$$

By using (2) we have

$$\psi(z+1) = \frac{d}{dz} \ln \Gamma(z+1) = \frac{d}{dz} \ln (z\Gamma(z)) = \frac{d}{dz} \ln z + \frac{d}{dz} \ln \Gamma(z) = \frac{1}{z} + \psi(z).$$

This proves the theorem. Iteration of (23) easily leads to

Theorem 13.

$$\psi(z+n) = \psi(z) + \frac{1}{z} + \frac{1}{z+1} + \dots + \frac{1}{z+n-1}, \quad n = 1, 2, 3, \dots \quad (24)$$

Another property of the digamma function is given by

Theorem 14.

$$\psi(z) - \psi(1-z) = -\frac{\pi}{\tan \pi z}, \quad z \neq 0, \pm 1, \pm 2, \dots \quad (25)$$

The proof of this theorem is based on (17). We have

$$\begin{aligned} \psi(z) - \psi(1-z) &= \frac{d}{dz} \ln \Gamma(z) + \frac{d}{dz} \ln \Gamma(1-z) = \frac{d}{dz} \ln (\Gamma(z)\Gamma(1-z)) \\ &= \frac{d}{dz} \ln \frac{\pi}{\sin \pi z} = \frac{\sin \pi z}{\pi} \cdot \frac{-\pi^2 \cos \pi z}{(\sin \pi z)^2} = -\frac{\pi}{\tan \pi z}. \end{aligned}$$

References

- [1] G.E. ANDREWS, R. ASKEY AND R. ROY, *Special Functions*. Encyclopedia of Mathematics and its Applications **71**, Cambridge University Press, 1999.