## Literature Study

## The Mild-Slope Equation and its Numerical Implementation

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TU Delft
December 8, 2011

## Outline

(1) The MSc Project

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(2) Geometry and Wave Motion

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4 Numerical Implementation

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4. Numerical Implementation
(5) Proposed Numerical Matrix Solver

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(7) Research Objectives

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- Outer iteration: Picard's method;
- Inner iteration: ILU-BiCGSTAB.
- Improve the undesired long computational time.


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## Geometry


$h$ Water height
$\zeta$ Elevation of the free surface
$a$ Absolute elevation, $a=|\zeta|$
$H$ Wave height, $H=2 A$
$k_{0}$ Wave number
$\omega$ Wave frequency
$L$ Wave length, $L=2 \pi / k_{0}$
$T$ Wave period, $T=2 \pi / \omega$

## Wave motion

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- The wave slope $\epsilon=\frac{2 \pi A}{L}$ is small;
- The wave motion is time harmonic;
- The changes in bottom topography are small;
- The surface tension can be neglected;
- The Coriolis effect can be neglected.


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Wave breaking - The process that causes large amounts of wave energy to be transformed into turbulent kinetic energy.

The energy dissipation due to wave breaking is described by the coefficient $W_{b}$ [1];

$$
W_{b}=\frac{2 \alpha}{T} Q_{b} \frac{H_{m}^{2}}{4 a^{2}}
$$

With

$$
\begin{aligned}
\alpha & \text { Adjustable constant; } \\
Q_{b} & \text { Fraction of breaking waves; } \\
H_{m} & \text { Maximal possible wave height; } \\
a & \text { Modulus of the free surface elevation, } a=|\tilde{\zeta}|
\end{aligned}
$$

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The energy dissipation due to bottom friction is described by the coefficient $W_{f}[4,22]$;

$$
W_{f}=\frac{8}{3 \pi} c_{f} \frac{a \omega^{3}}{\sinh ^{3}\left(k_{0} h\right)}
$$

With
$c_{f}$ Bottom friction coefficient.

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## The Mild-Slope equation

The Mild-Slope equation, with energy dissipation included, is given by [4]

$$
\begin{equation*}
\nabla \cdot\left(\frac{n_{0}}{k_{0}^{2}} \nabla \tilde{\zeta}\right)+\left(n_{0}-\frac{i W}{\omega}\right) \tilde{\zeta}=0 \tag{1}
\end{equation*}
$$

with
$W$ The energy dissipation term $W=W_{f}+W_{b}$;
$n_{0}$ A constant $n_{0}=\frac{1}{2}\left(1+\frac{2 k_{0} h}{\sinh \left(2 k_{0} h\right)}\right)$;
$\nabla=\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}\right)^{T}$.

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Non-linearity:

$$
W \tilde{\zeta}=\left(\frac{8}{3 \pi} c_{f} \frac{|\tilde{\zeta}| \omega^{3}}{\sinh ^{3}\left(k_{0} h\right)}+\frac{2 \alpha}{T} Q_{b} \frac{H_{m} 2}{4|\tilde{\zeta}|^{2}}\right) \tilde{\zeta}
$$

## The Mild-Slope equation

There are two types of boundaries, i.e.

- The open boundary with an incoming and an outgoing wave.
- The closed boundary where partial reflection due to interaction with the boundary occurs.


## The Mild-Slope equation

Boundary conditions

The condition for the open boundary $\left(\Gamma_{1}\right)$;

$$
\frac{\partial \tilde{\zeta}}{\partial n}=-p \tilde{\zeta}_{i n}\left(\boldsymbol{e}_{i n} \cdot \boldsymbol{n}\right)-p\left(\tilde{\zeta}-\tilde{\zeta}_{i n}\right)+\frac{1}{2 p}\left(\frac{\partial^{2} \tilde{\zeta}}{\partial s^{2}}+\frac{\partial^{2} \tilde{\zeta}_{i n}}{\partial s^{2}}\right)
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The condition for the closed boundary $\left(\Gamma_{2}\right)$;

$$
\frac{\partial \tilde{\zeta}}{\partial n}=-\left(\frac{1-R}{1+R}\right)\left\{p \tilde{\zeta}-\frac{1}{2 p} \frac{\partial^{2} \tilde{\zeta}}{\partial s^{2}}\right\}
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$$

With
$R$ the reflection coefficient, $0 \leq R \leq 1$;
$p$ The modified wave number $p=i k_{0} \sqrt{1-\frac{i W}{\omega n_{0}}}$.

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## Currently used Numerical Methods

- Ritz-Galerkin Finite Element Method
- Bi-CGSTAB
- Incomplete LU factorization
- Picard iteration


## Ritz-Galerkin Finite Element Method

## Weak formulation

The weak formulation of the Mild-Slope equation is given by

$$
\begin{aligned}
& \int_{\Omega}\left\{\left(n_{0}-\frac{i W}{\omega}\right) \tilde{\zeta} \eta-\frac{n_{0}}{k_{0}^{2}} \nabla \tilde{\zeta} \cdot \nabla \eta\right\} d \Omega \\
& -\int_{\Gamma_{2}} \frac{n_{0}}{k_{0}^{2}}\left(\frac{1-R}{1+R}\right)\left\{p \tilde{\zeta} \eta+\frac{1}{2 p} \frac{\partial \tilde{\zeta}}{\partial s} \frac{\partial \eta}{\partial s}\right\} d \Gamma \\
& -\int_{\Gamma_{1}} \frac{n_{0}}{k_{0}^{2}}\left\{p \tilde{\zeta} \eta+\frac{1}{2 p} \frac{\partial \tilde{\zeta}}{\partial s} \frac{\partial \eta}{\partial s}\right\} d \Gamma \\
& =\int_{\Gamma_{1}} \frac{n_{0}}{k_{0}^{2}}\left\{p \tilde{\zeta}_{\text {in }}\left(\boldsymbol{e}_{\text {in }} \cdot \boldsymbol{n}\right) \eta-p \tilde{\zeta}_{\text {in }} \eta-\frac{1}{2 p}\left(\frac{\partial \tilde{\zeta}_{\text {in }}}{\partial s} \frac{\partial \eta}{\partial s}\right)\right\} d \Gamma .
\end{aligned}
$$

## Ritz-Galerkin Finite Element Method

## Resulting matrix

The domain is divided into triangles, with piecewise linear basis functions. This results in the following system

$$
\boldsymbol{A} \boldsymbol{x}=\boldsymbol{y}
$$

$\boldsymbol{A} \in \mathbb{C}^{N \times N}, \boldsymbol{x} \in \mathbb{C}^{N}$ and $\boldsymbol{y} \in \mathbb{C}^{N}$. With

$$
\boldsymbol{A}=\left(-\boldsymbol{L}-\boldsymbol{C}+z_{1} \boldsymbol{M}\right) .
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$$
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& \quad+\int_{\Gamma_{1}} \frac{n_{0}}{k_{0}^{2}}\left\{p \tilde{\zeta} \eta+\frac{1}{2 p} \frac{\partial \tilde{\zeta}}{\partial s} \frac{\partial \eta}{\partial s}\right\} d \Gamma \Rightarrow C
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With

$$
\int_{\Omega} \tilde{\zeta} \eta d \Omega \quad \Rightarrow \quad \boldsymbol{M} \quad \text { and } \quad z_{1}=\left(n_{0}-\frac{i W}{\omega}\right)
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## Numerical Matrix Solver

Bi-CGSTAB

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- Bi-CGSTAB is a Krylov subspace method, the Krylov subspace of dimension $m$ is given by

$$
\mathcal{K}_{m}\left(\boldsymbol{A} ; \boldsymbol{r}_{0}\right)=\operatorname{span}\left\{\boldsymbol{r}_{0}, \boldsymbol{A} \boldsymbol{r}_{0}, \ldots, \boldsymbol{A}^{m-1} \boldsymbol{r}_{0}\right\}
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- The residual of $\mathrm{Bi}-\mathrm{CGSTAB}$ can be written as

$$
\boldsymbol{r}_{i}^{B i-C G S T A B}=Q_{i}(\boldsymbol{A}) P_{i}(\boldsymbol{A}) \boldsymbol{r}_{0}
$$

with

$$
Q_{i}(\boldsymbol{A})=\left(I-\omega_{1} \boldsymbol{A}\right)\left(I-\omega_{2} \boldsymbol{A}\right) \ldots\left(I-\omega_{i} \boldsymbol{A}\right)
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## Numerical Matrix Solver

Incomplete LU factorization

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- ILU(0) is based on
- Matrices $\boldsymbol{L}$ and $\boldsymbol{U}$ have the same zero-pattern as $\boldsymbol{A}$, i.e. $u_{i, j}=l_{i, j}=0$ if $a_{i, j}=0$ and if $a_{i, j} \neq 0$ then $u_{i, j} \neq 0$ and $l_{i, j} \neq 0$.
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- The LU decomposition proposed by A. van der Ploeg [16] is used in HARES.


## Picard iteration

The following algorithm is used for the Picard iteration:
$W^{0}=0$; initial value for the dissipation term
for $\quad i=1,2, \ldots$
Solve $\boldsymbol{x}^{i}$ from $\boldsymbol{A}\left(W^{i-1}\right) \boldsymbol{x}^{i}=\boldsymbol{b}$
$W^{i}=W\left(\boldsymbol{x}^{i}\right)$
end

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- Generate residuals $\boldsymbol{r}_{n}$ that are in the subspace $\mathcal{G}_{j}$ with decreasing dimension;

$$
\mathcal{G}_{j}=\left(\boldsymbol{I}-\omega_{j} \boldsymbol{A}\right)\left(\mathcal{G}_{j-1} \cap \boldsymbol{P}^{\perp}\right),
$$

with $\mathcal{G}_{0}=\mathcal{K}^{N}\left(\boldsymbol{A} ; \boldsymbol{v}_{0}\right)$ and $\boldsymbol{P} \in \mathbb{C}^{N \times s}$.

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- Based on the IDR theorem [13], which states that
(i) $\mathcal{G}_{j} \subset \mathcal{G}_{j-1}$ for all $\mathcal{G}_{j-1} \neq\{\mathbf{0}\}, j>0$,
(ii) $\mathcal{G}_{j}=\{\mathbf{0}\}$ for some $j \leq N$


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- Residuals are obtained by

$$
\boldsymbol{r}_{n+1}=\left(\boldsymbol{I}-\omega_{j+1} \boldsymbol{A}\right) \boldsymbol{v}_{n} \quad \text { with } \quad \boldsymbol{v}_{n} \in \mathcal{G}_{j} \cap \boldsymbol{P}^{\perp}
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- Requires at most $N+\frac{N}{s}$ matrix-vector multiplications.


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Freedom in the $\operatorname{IDR}(s)$ algorithm

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- Choosing $P$.


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Currently orthogonalized random vectors are used.

## Proposed Numerical Matrix Solver IDR(s)

Freedom in the $\operatorname{IDR}(s)$ algorithm

- Choosing $\boldsymbol{P}$.

$$
\mathcal{G}_{j}=\left(\boldsymbol{I}-\omega_{j} \boldsymbol{A}\right)\left(\mathcal{G}_{j-1} \cap \boldsymbol{P}^{\perp}\right)
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- Choosing $\omega_{j}$. For each residual in $\mathcal{G}_{j}$ the same $\omega_{j}$ is needed, currently based on a strategy proposed by Sleijpen and Van der Vorst. [11]


## Proposed Numerical Matrix Solver IDR(s)

Freedom in the $\operatorname{IDR}(s)$ algorithm

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(i) Can be done using a simple Krylov method;


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(i) Can be done using a simple Krylov method;
(ii) Can be chosen freely as long as $\mathcal{G}_{0}$ is the complete Krylov subspace $\mathcal{K}^{N}$.


## Proposed Numerical Matrix Solver

## Shifted Laplace preconditioner

- The shifted Laplace preconditioner [6], used in our experiments, is given by

$$
\boldsymbol{K}=\left(-\boldsymbol{L}-\boldsymbol{C}-i\left|z_{1}\right| \boldsymbol{M}\right),
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with $z_{1}=\left(n_{0}-\frac{i W}{\omega}\right)$.

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If $L$ and $C$ are symmetric positive semidefinite real matrices and $\boldsymbol{M}$ symmetric positive definite real matrix, that the eigenvalues of the preconditioned system lie inside or on a circle. [19]

## Outline

## (1) The MSc Project

(2) Geometry and Wave Motion
(3) The Mild-Slope Equation
(4) Numerical Implementation
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## Numerical Experiments

Problem dimension \& used matrix-solvers

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The following test problem is considered:

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- 126.504 internal triangular elements;
- 2.213 boundary line segments;
- 63.253 unknowns;


## Numerical Experiments

Problem dimension \& used matrix-solvers

The following test problem is considered:

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- ILU(0), ILU(0)-Shifted Laplace.


## Numerical Experiments

## Number of matrix-vector products


(c) Number of matvecs for the ILU(0) preconditioned system

(d) Number of matvecs for the ILU(0)shifted laplace preconditioned system

## Numerical Experiments

## Computational time

|  | Numerical method |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | Bi-CGSTAB | CG-S | IDR(2) | IDR(4) |
| ILU(0) | $3.0192 \cdot 10^{3}$ | $4.1053 \cdot 10^{3}$ | $1.2303 \cdot 10^{3}$ | $1.2863 \cdot 10^{3}$ |
| ILU(0)-SL | $1.8525 \cdot 10^{3}$ | $2.6634 \cdot 10^{3}$ | $0.7361 \cdot 10^{3}$ | $0.9741 \cdot 10^{3}$ |

Table: CPU time until the whole process is completed.

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Table: CPU time until the whole process is completed.

Using $\operatorname{IDR}(s)$ preconditioned with the incomplete LU factorization of the shifted Laplace matrix speeds the computational time up with a factor three.

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## Research objectives

## Improvement of HARES :

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- The non-linear part;


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(iii) Total error.


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Improvement of HARES :

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Theoretical research:

## Research objectives

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Theoretical research:

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(i) Analysis in [19] not identical for Mild-Slope equation;
(ii) Determine an optimal shift when possible;
(iii) How should we approximate the shifted Laplace preconditioner.


## Research objectives

Improvement of HARES :

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When the spectrum is a circle we might be able to determine the Ritz-values such that the polynomial

$$
Q_{j}(\boldsymbol{A})=\left(\boldsymbol{I}-\omega_{1} \boldsymbol{A}\right) \ldots\left(\boldsymbol{I}-\omega_{j} \boldsymbol{A}\right)
$$

has a minimal maximum on the spectrum.

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Improvement of HARES :

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- Implementation into FORTRAN.

Theoretical research:

- Spectral analysis for preconditioner of the Mild-Slope equation;
- Choosing the coefficients $\omega_{j}$ based on Ritz-values;
- Can smartly building $\mathcal{G}_{0}$ lead to convergence speed up.


## Questions

## QUESTIONS?

## Figure diffraction \& reflection



Figure: The harbour of Scheveningen. The effects of diffraction - reflection visible

## Figure refraction \& shoaling


(a) Refraction

(Plummer et al., 2001)
(b) Shoaling

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